

# Kagome Networks and Frustration

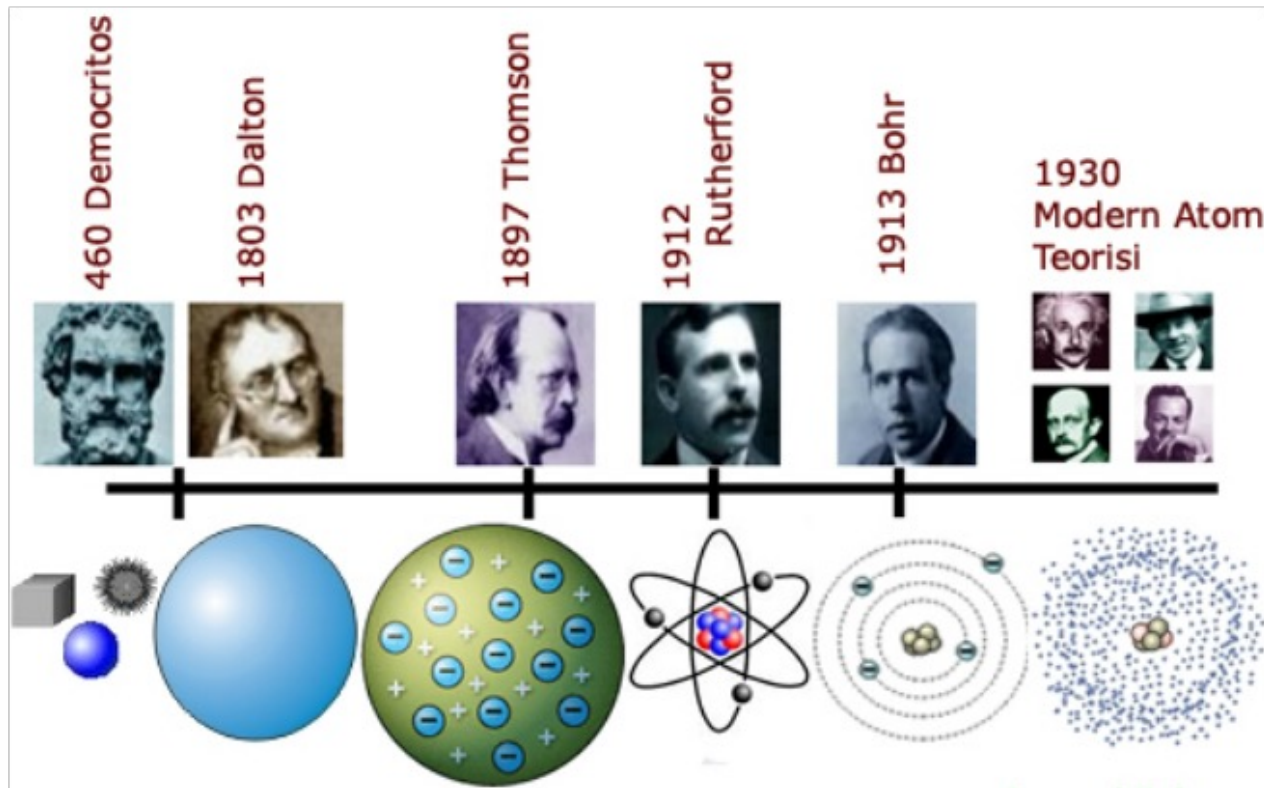
Andrea Capa Salinas  
Matr 286G  
05/31/2022

# Emergent Phenomena



An entity exhibits properties its parts do not possess on their own

# Reductionist Approach



<https://saintschemistry10.weebly.com/history-of-the-atom.html>



# Reductionist Approach

## Standard Model of Elementary Particles

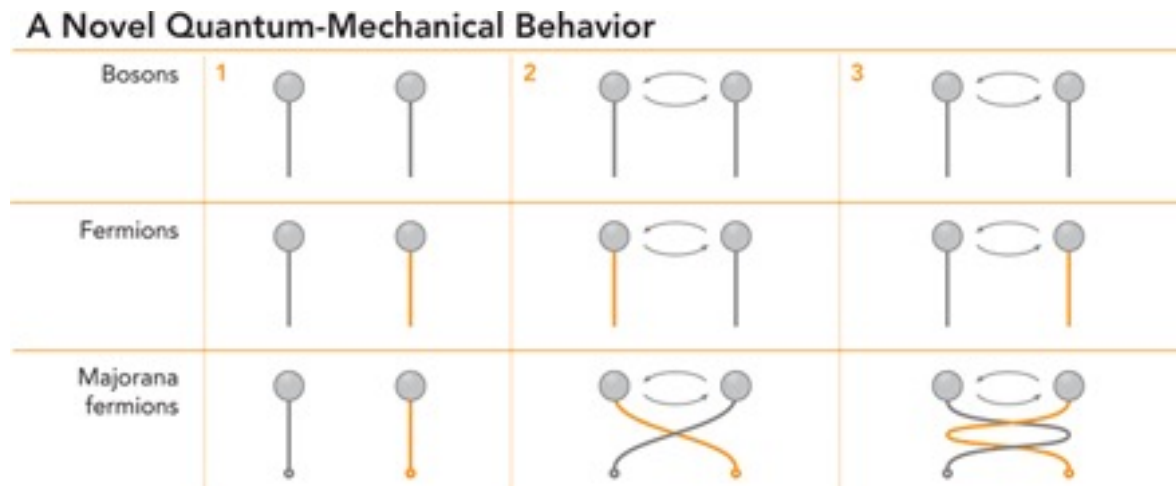
		three generations of matter (elementary fermions)			three generations of antimatter (elementary antifermions)			interactions / force carriers (elementary bosons)	
		I	II	III	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	0	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0	0
	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b><math>\bar{u}</math></b> antiup	<b><math>\bar{c}</math></b> anticharm	<b><math>\bar{t}</math></b> antitop	<b>g</b> gluon		<b>H</b> higgs
<b>QUARKS</b>	$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	0	
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}$	0	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0	
	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom	<b><math>\bar{d}</math></b> antidown	<b><math>\bar{s}</math></b> antistrange	<b><math>\bar{b}</math></b> antibottom	<b><math>\gamma</math></b> photon		
	$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 91.19 \text{ GeV}/c^2$	0	
	-1	-1	-1	1	1	1	0	1	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	1	
	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b><math>e^+</math></b> positron	<b><math>\bar{\mu}</math></b> antimuon	<b><math>\bar{\tau}</math></b> antitau	<b>Z</b> Z <sup>0</sup> boson		
<b>LEPTONS</b>	$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$\approx 80.39 \text{ GeV}/c^2$	1	
	0	0	0	0	0	0	1	1	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	1	
	<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino	<b><math>\bar{\nu}_e</math></b> electron antineutrino	<b><math>\bar{\nu}_\mu</math></b> muon antineutrino	<b><math>\bar{\nu}_\tau</math></b> tau antineutrino	<b>W<sup>+</sup></b> W <sup>+</sup> boson		<b>W<sup>-</sup></b> W <sup>-</sup> boson

**GAUGE BOSONS**  
VECTOR BOSONS

**SCALAR BOSONS**

# Emergent Phenomena in Kagome Networks

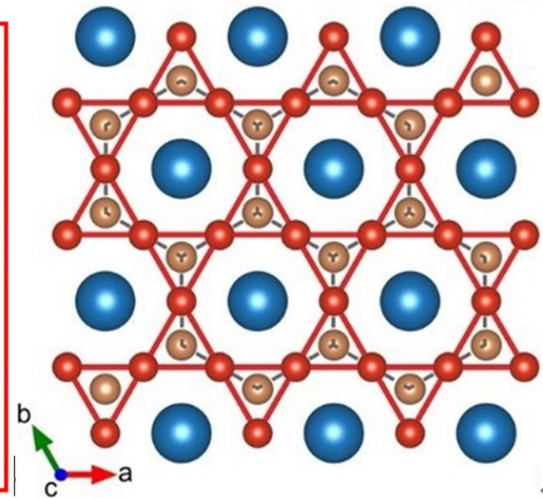
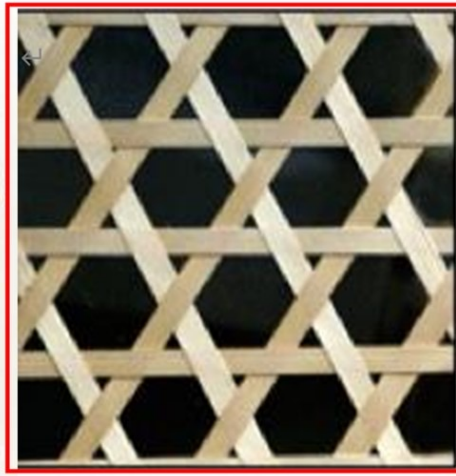
- Fractional Quantum Hall Effect
- Quantum Spin Liquid States
- Spinons
- Majorana fermions



Service R. F. (2011). Search for Majorana fermions nearing success at last?. *Science*. 332(6026), 193–195.

# Kagome Networks

- Intersecting webs of “corner-sharing triangles”
- Named after patterned bamboo basket
  - Bamboo-basket (kago) woven pattern (me)



[https://english.cas.cn/newsroom/research\\_news/phys/202112/t20211217\\_294587.shtml](https://english.cas.cn/newsroom/research_news/phys/202112/t20211217_294587.shtml)

# Kagome Networks

- Lars Onsager (1944) paper on square lattice ising model
- Kodi Husimi and Itiro Syôzi (1950) simplified Onsagers solutions and solved for honeycomb and triangular
- Syôzi (1951) studied Kagome lattice

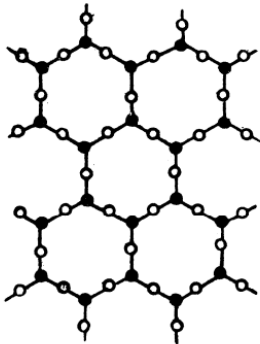


Fig. 2. Decorated Honeycomb Lattice

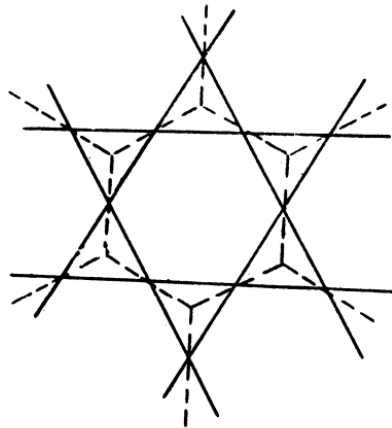


Fig. 3. Star-Triangle Transformation

I. Syôzi, Prog. Theor. Phys. 6, 306 (1951).

Mamoru Mekata , "Kagome: The Story of the Basketweave Lattice", Physics Today 56, 12-13 (2003)

Progress of Theoretical Physics, Vol. VI, No. 3, May~June, 1951.

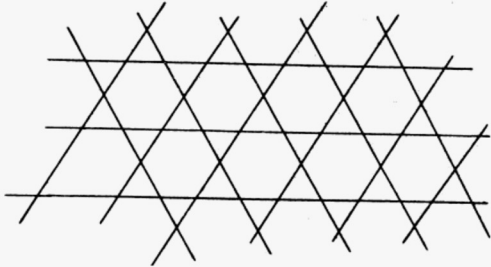
**Statistics of Kagomé Lattice**

Itiro Syôzi  
*Department of Physics, Osaka University*

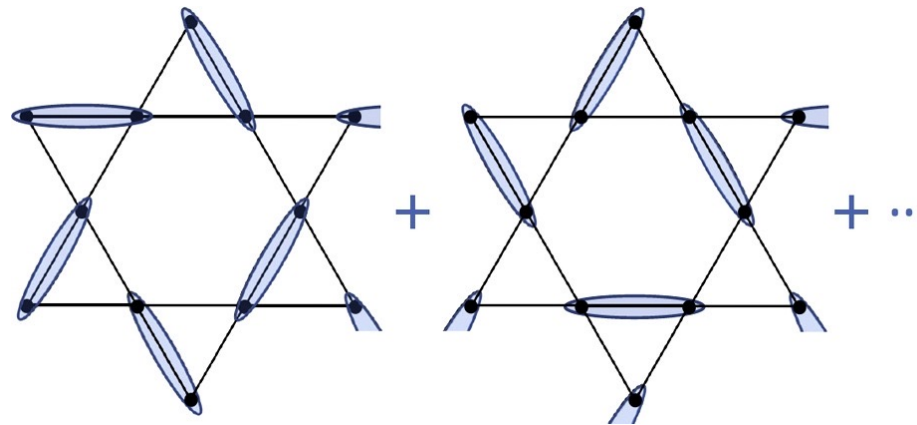
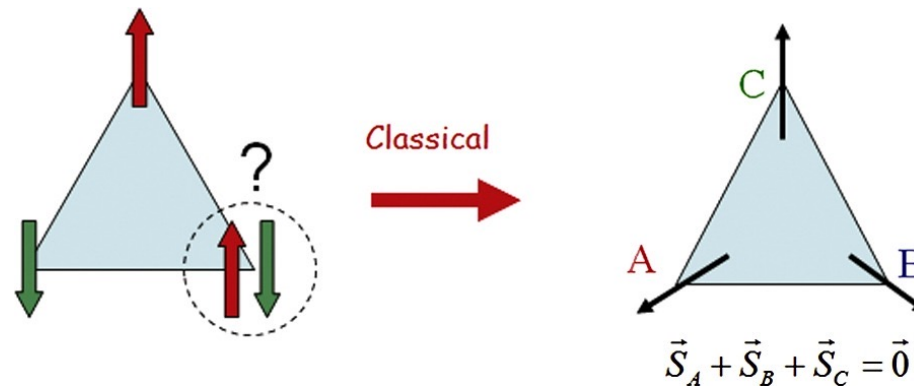
(Received February 27, 1951)

The transition temperature of the kagomé lattice with  $Z=4$  is obtained and compared with that of the square lattice.

After the work of Onsager,<sup>1)</sup> who solved exactly the problem of Ising model for the case of plane square lattice, the same problems for the honeycomb and triangular lattice were treated by several authors.<sup>2)</sup> Other than these three types of lattices, there is left a lattice, called in Japanese kagomé (woven bamboo pattern), which consists exclusively of equivalent lattice points and equivalent bonds. Since the number of nearest neighbors of a lattice point is as many as in the square lattice, namely four, it is interesting to verify the natural conjecture that the curie point, in general, is determined solely by the relation  $\chi^2 H = \sec \pi/Z$  established by Onsager for the three types of lattices.

A diagram of a Kagomé lattice, which is a tiling of the plane by triangles. The lattice is shown in a perspective view, with some edges and vertices highlighted in black.

# Frustration in Kagome Network



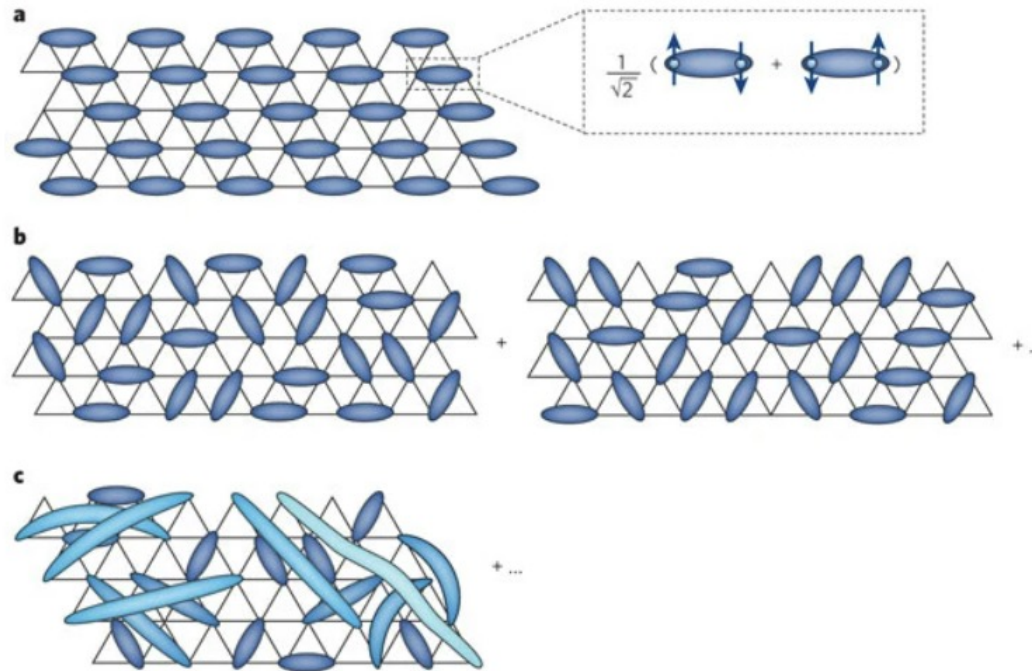
Mendels, P., & Bert, F. (2016). *Comptes Rendus Physique*, 17(3-4), 455-470.



# Resonant Valence Bond Theory

- Valence bonds are allowed to undergo quantum mechanical fluctuations
- GS is a superposition of different partitionings of spins into valence bonds

Figure 3: Valence-bond states of frustrated antiferromagnets.



Mendels, P., & Bert, F. (2016). *Comptes Rendus Physique*, 17(3-4), 455-470.

# First Kagome Candidate

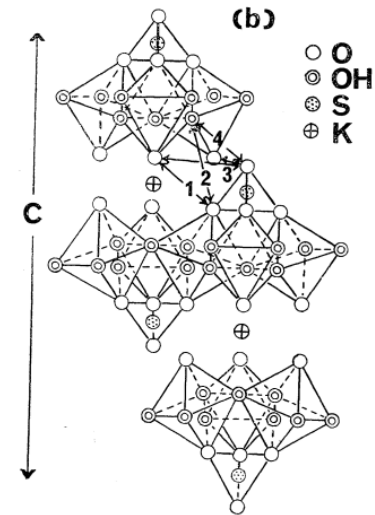
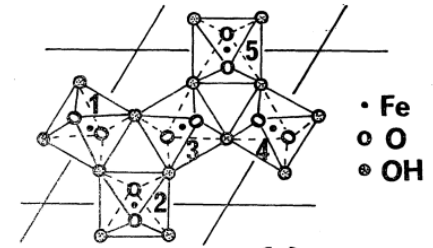
## On the Spin Arrangement in “Kagome” Lattice of Antiferromagnetic $\text{KFe}_3(\text{OH})_6(\text{SO}_4)_2$

Mikio TAKANO, Teruya SHINJO and Toshio TAKADA

*Institute for Chemical Research, Kyoto University, Uji, Kyoto-fu*

(Received August 18, 1970)

Experimental results showed that  $\text{KFe}_3(\text{OH})_6(\text{SO}_4)_2$  is antiferromagnetic below  $60^\circ\text{K}$  and in each crystal  $c$ -plane, Fe ions form a compensated anti-ferromagnet. The crystal structure of Fe ions in the  $c$ -plane is a “kagome” lattice and it has been known that collinear antiferromagnetic spin arrangements are not stable in the kagome lattice. An estimation of the spin structure was attempted with the analysis of the Mössbauer spectrum and a kind of triangular configuration was suggested.



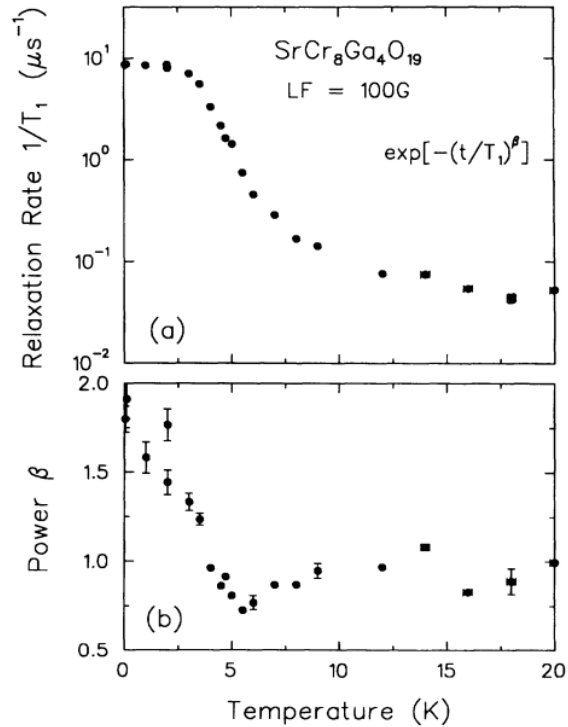
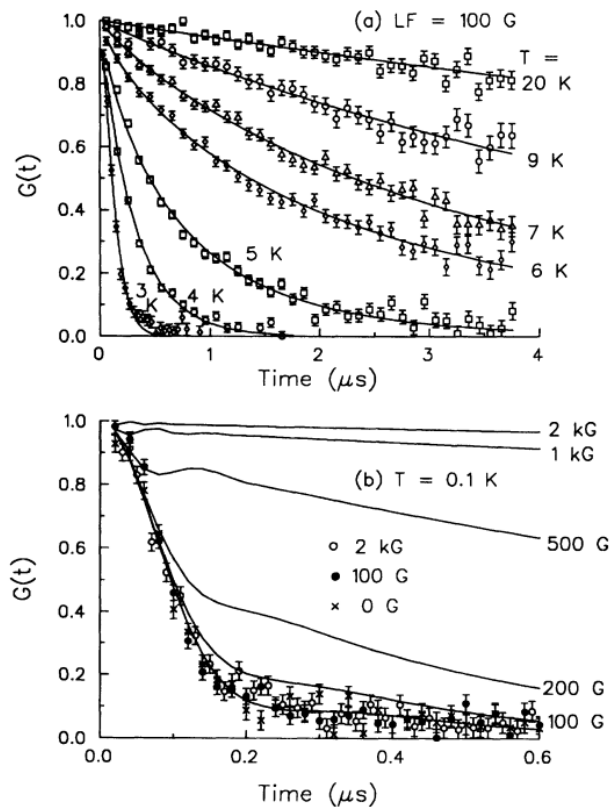
## Jarosite



M. Takano, T. Shinjo, M. Kiyama, T. Takada, J. Phys. Soc. Jpn. 25, 902 (1968)

<https://underthescopeminerals.tumblr.com/post/178252086249/jarosite-kfe3-3so42oh6-localityla>

# Kagome System $\text{SrCr}_8\text{Ga}_4\text{O}_{19}$

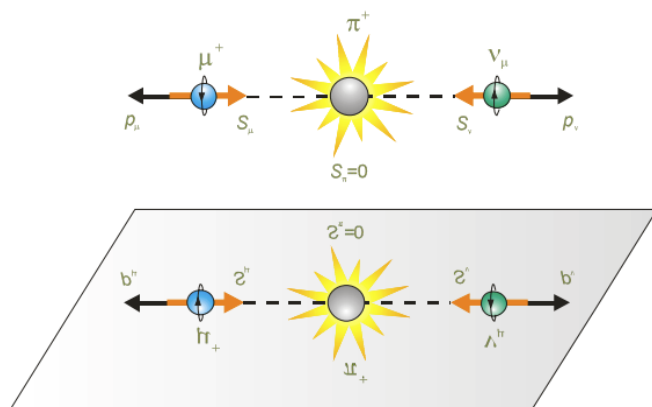
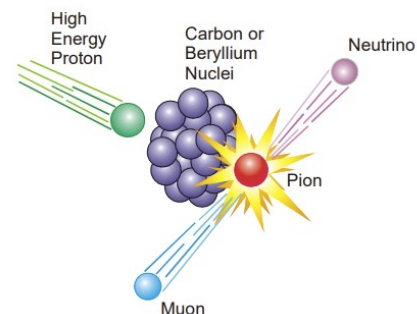


Uemura, Y. J., Keren, A., Kojima, K., Le, L. P., Luke, G. M., Wu, W. D., ... & Kakurai, K. Phys Rev Lett. 73(24), 3306. (1994)

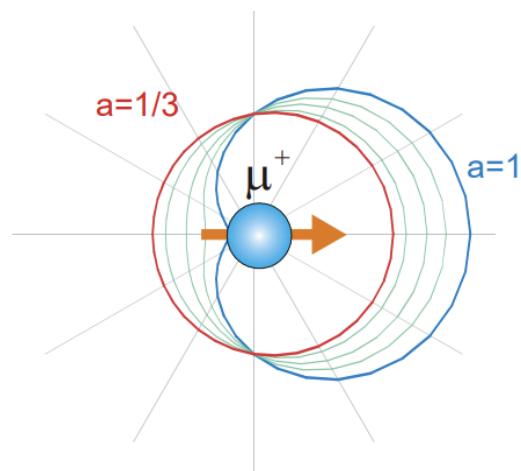
# $\mu$ SR Muon Spin Relaxation Technique

$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

$$\pi^- \rightarrow \mu^- + \bar{\nu}_\mu$$



Parity-violating collinear decay of a pion  $\pi^+$  at rest into a muon  $\mu^+$  and a muonic neutrino  $\nu_\mu$ .



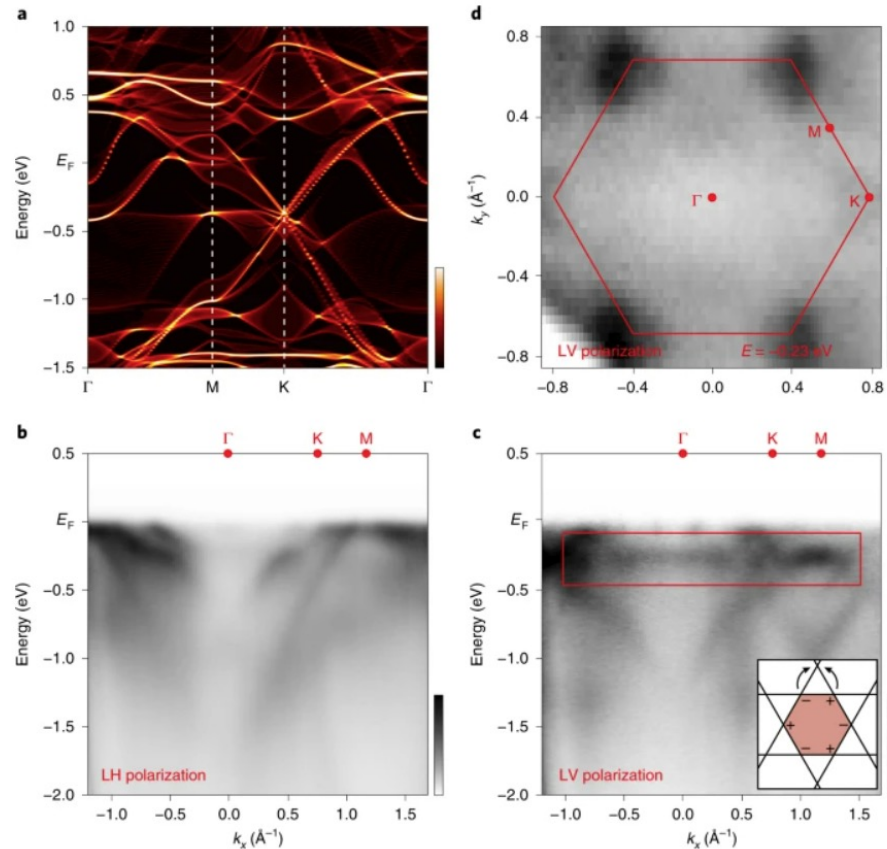
Angular distribution of the positrons from the muon decay:  $W(E,\theta) = 1 + a(E)\cos(\theta)$ . When all positron energies  $E$  are sampled with equal probability the asymmetry parameter has the

<http://www.chem.ubc.ca/sites/default/files/users/dgf/musrbrochure.pdf>



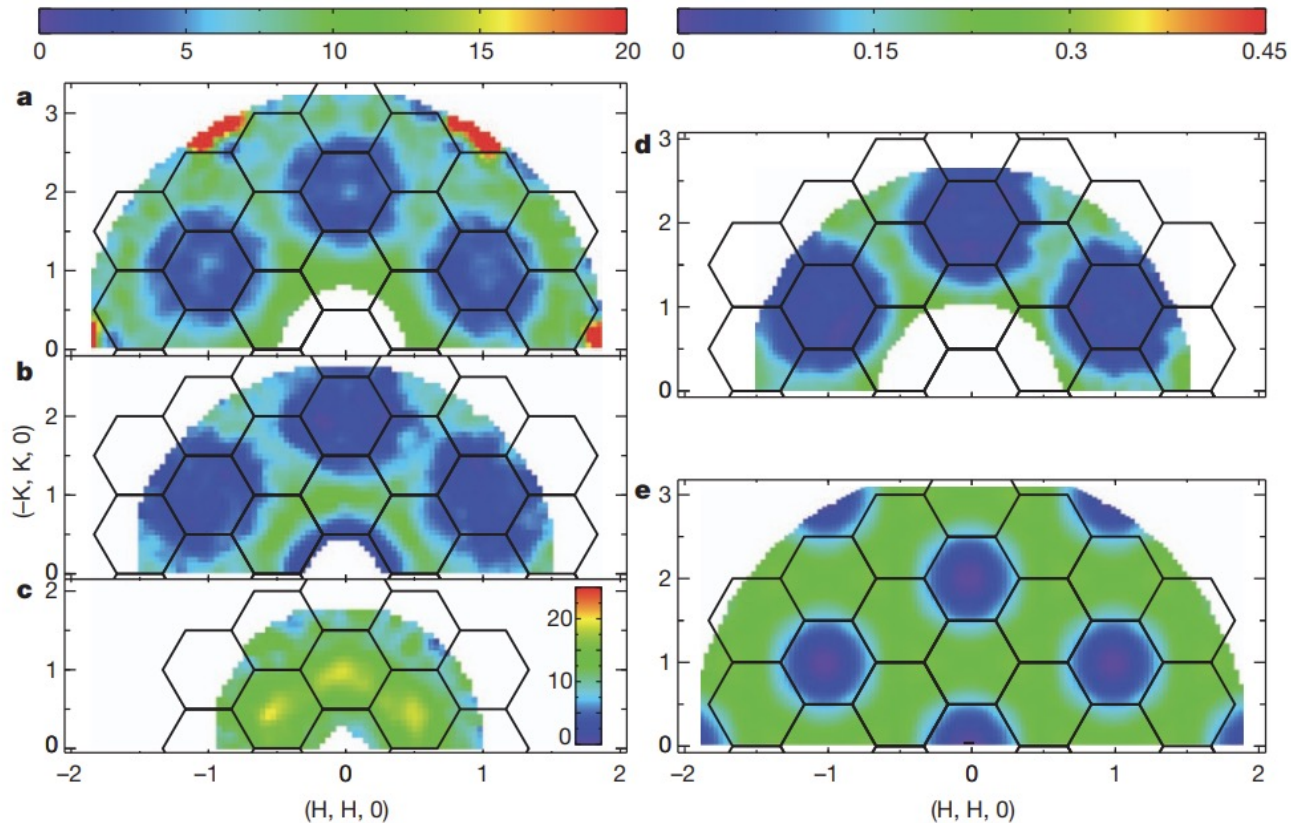
# Kagome FeSn

Fig. 4: Signature of flat bands in FeSn.



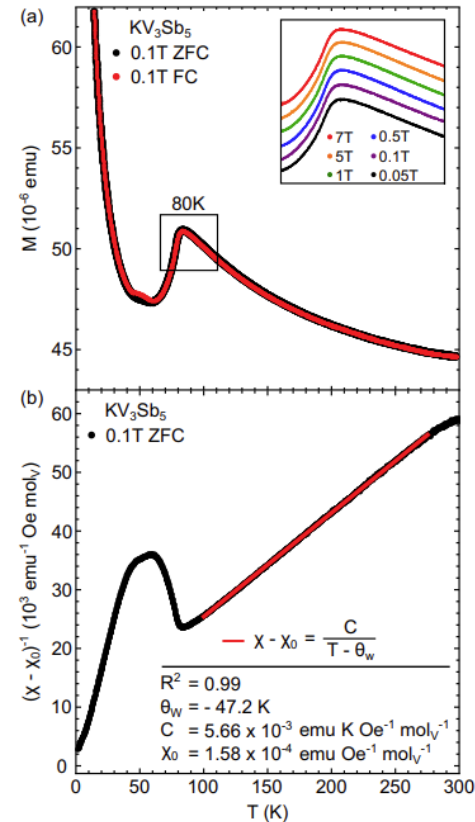
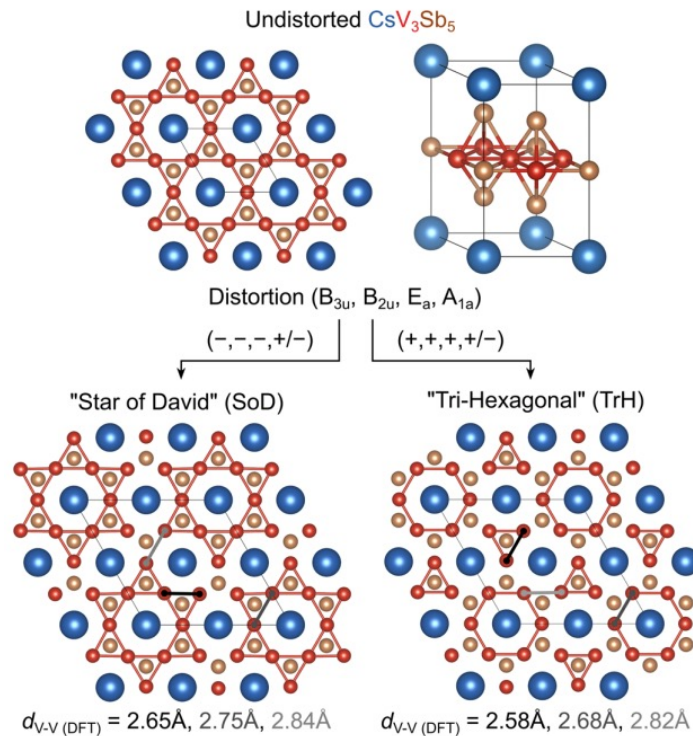
Kang, M., Ye, L., Fang, S. et al. Dirac fermions and flat bands in the ideal kagome metal FeSn. Nat. Mater. 19, 163–169 (2020).

# $\text{ZnCu}_3(\text{OD})_6\text{Cl}_2$ (herbertsmithite)



Han, T. H., Helton, J. S., Chu, S., Nocera, D. G., Rodriguez-Rivera, J. A., Broholm, C., & Lee, Y. S. (2012). Fractionalized excitations in the spin-liquid state of a kagome-lattice antiferromagnet. *Nature*, 492(7429), 406–410.

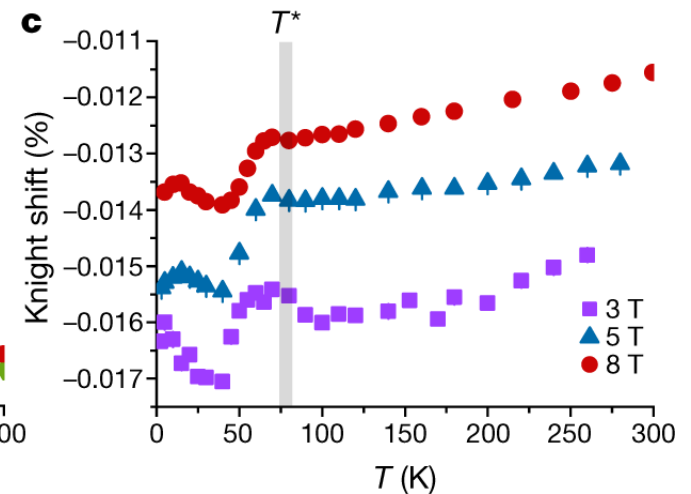
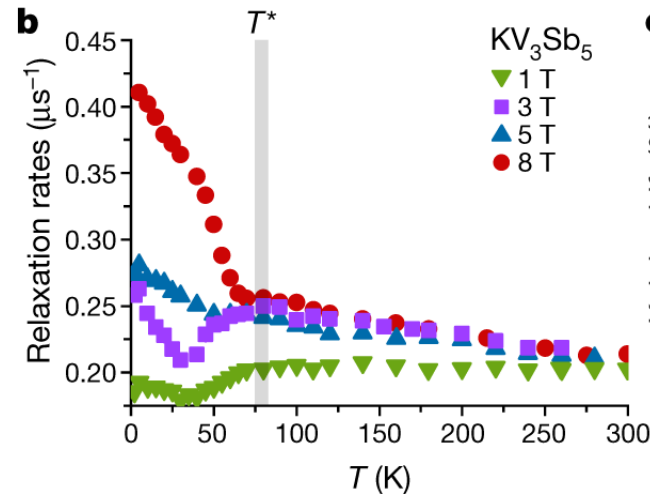
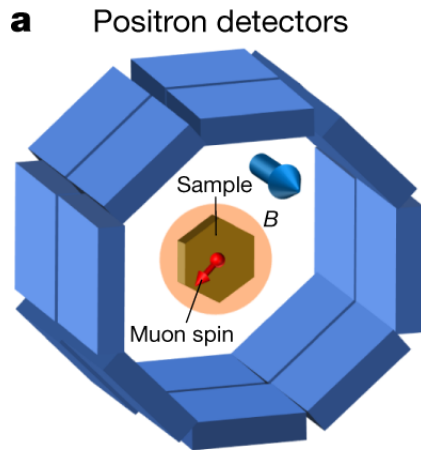
# $AV_3Sb_5$ (A=K,Rb,Cs)



Ortiz, B. R; Teicher, S.; Kautzsch, L.; Sarte, P. M; Ratcliff, N.; Harter, J., et al. Physical Review X, 11(4). (2021).

Ortiz, Brenden R., Gomes, LÍdia C., Morey, Jennifer R., Winiarski, Michal, Bordelon, Mitchell, Mangum, John S., Oswald, Iain W. H., Rodriguez-Rivera, Jose A., Neilson, James R., Wilson, Stephen D., Ertekin, Elif, McQueen, Tyrel M., & Toberer, Eric S. (2019)

# $AV_3Sb_5$ (A=K,Rb,Cs)



Mielke, C., Das, D., Yin, JX. et al. Time-reversal symmetry-breaking charge order in a kagome superconductor. Nature 602, 245–250 (2022).





Thank You!